

## NOVAE

# Fermi establishes classical novae as a distinct class of gamma-ray sources

The Fermi-LAT Collaboration\*†

A classical nova results from runaway thermonuclear explosions on the surface of a white dwarf that accretes matter from a low-mass main-sequence stellar companion. In 2012 and 2013, three novae were detected in  $\gamma$  rays and stood in contrast to the first  $\gamma$ -ray-detected nova V407 Cygni 2010, which belongs to a rare class of symbiotic binary systems. Despite likely differences in the compositions and masses of their white dwarf progenitors, the three classical novae are similarly characterized as soft-spectrum transient  $\gamma$ -ray sources detected over 2- to 3-week durations. The  $\gamma$ -ray detections point to unexpected high-energy particle acceleration processes linked to the mass ejection from thermonuclear explosions in an unanticipated class of Galactic  $\gamma$ -ray sources.

The Fermi-LAT (Large Area Telescope) (1), launched in 2008, continuously scans the sky in  $\gamma$  rays, thus enabling searches for transient sources. When a nova explodes in a symbiotic binary system, the ejecta from the white dwarf surface expand within the circumstellar wind of the red giant companion, and high-energy particles can be accelerated in a blast wave driven in the high-density environment (2) so that variable  $\gamma$ -ray emission can be produced,

as was detected at  $>100$ -MeV energies by the LAT in V407 Cygni 2010 (V407 Cyg) (3). In a classical nova, by contrast, the ejecta quickly expand beyond the confines of the compact binary into a much lower-density environment. High-energy particle acceleration could therefore be related to a bow shock driven by the ejecta in the interstellar medium or to turbulence and eventually weaker internal shocks formed in the inhomogeneous ejecta itself. The contribution of such expanding nova shells to cosmic-ray acceleration has been considered (4), but no predictions have so far been made for  $>100$ -MeV  $\gamma$  rays. The classical novae (or simply “novae” where appropriate) detected by the LAT with statistical significance

of 12 to 20  $\sigma$  (Table 1 and Fig. 1)—V959 Monocerotis 2012 (V959 Mon), V1324 Scorpii 2012 (V1324 Sco), and V339 Delphini 2013 (V339 Del)—were unanticipated. These observed  $\gamma$  rays have higher energies than nuclear line emission by radioactive decay at  $\sim$  MeV energies that remain undetected in individual novae (5) and  $\leq$  0.1-MeV emission detected in isolated cases (6).

V959 Mon was detected as a transient  $\gamma$ -ray source in June 2012 by the LAT while close ( $\sim 20^\circ$  separation) to the Sun (7) and then optically in August (8). Ultraviolet spectroscopy revealed an oxygen-neon nova (9), recognized as the class with the most massive white dwarfs ( $\geq 1.1 M_\odot$ ) with massive ( $\geq 8 M_\odot$ ) progenitors [e.g., (10)]. The expected peak visual magnitude of  $\sim 5$  would have been observable with the naked eye  $\sim 50$  days earlier, when the  $\gamma$ -ray transient was detected (9). V339 Del (11) was detected in August 2013 in a LAT-pointed observation triggered by its high optical brightness [4.3 magnitude at peak (12, 13)]. Optical spectra of V339 Del suggest a carbon-oxygen nova (14), which are more common than the oxygen-neon types, with less massive white dwarfs evolved from  $\leq 8 M_\odot$  main-sequence progenitors. Optical brightening of V1324 Sco was detected in May 2012 (15) and found in LAT  $\gamma$ -ray data from June (16). Although the type for V1324 Sco is currently unclear, its optical spectroscopic evolution at early times (15) did not resemble oxygen-neon novae at similar stages. We take this to indicate that it is likely a carbon-oxygen type.

The LAT data (13) for the three classical novae are discussed together with an updated analysis of the originally detected symbiotic nova V407

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**Table 1. Summary of the four novae.** Tabulated are optical peak magnitudes and adopted distances from (19) for V407 Cyg, estimate of  $\sim 4$ - to 5-kpc V1324 Sco based on the maximum magnitude rate of decline relation (17) [notwithstanding the large uncertainties in this method (29)], (9) for V959 Mon (scaled from V1974 Cyg 1992), and (30) for V339 Del (scaled from OS Andromedae 1986), and observed dates of the optical peaks [unfiltered from (3), V-band, adopted, and visual magnitudes, respectively]. Positions in J2000.0 equinox (right ascension, RA; declination, Decl.), Galactic longitude ( $l$ ) and latitude ( $b$ ), 95% confidence localization error

radius, and offset between the LAT and optical positions in units of degrees. Adopted start dates  $t_s$  (13) are given in Gregorian dates and modified Julian days (MJD). The  $\gamma$ -ray luminosities  $L_\gamma$  and total emitted energies were estimated with the average fluxes from the power-law fits of the  $>100$ -MeV LAT spectra integrated up to 10 GeV and durations from  $t_s$  up to the last  $>2 \sigma$  daily bin LAT detection. For V339 Del, the  $\gamma$  rays were detected for 25 days in 1-day bins (Fig. 2), but there was a hint of a detection 2 days earlier on the day of the optical peak in 0.5-day binned data (13), leading to a 27-day duration.

Nova	V407 Cyg 2010	V1324 Sco 2012	V959 Mon 2012	V339 Del 2013
Distance (kpc)	2.7	4.5	3.6	4.2
Peak magnitude	6.9	10.0	5*	4.3
Peak date	10.80 Mar 2010	19.96 Jun 2012	—	16.50 Aug 2013
Optical RA, Decl.	315.5409°, +45.7758°	267.7246°, −32.6224°	99.9108°, +5.8980°	305.8792°, +20.7681°
Optical $l, b$	86.9826°, −0.4820°	357.4255°, −2.8723°	206.3406°, +0.0754°	62.2003°, −9.4234°
LAT RA, Decl.	315.57°, +45.75°	267.72°, −32.69°	99.98°, +5.86°	305.91°, +20.78°
Optical-LAT offset	0.03°	0.07°	0.08°	0.03°
LAT error radius (95%)	0.08°	0.09°	0.18°	0.12°
$t_s$ (date)	10 Mar 2010	15 Jun 2012	19 Jun 2012	16 Aug 2013
$t_s$ (MJD)	55265	56093	56097	56520
Duration (days)	22	17	22	27
$L_\gamma$ ( $10^{35}$ erg s $^{-1}$ )	3.2	8.6	3.7	2.6
Total energy ( $10^{41}$ erg)	6.1	13	7.1	6.0

\*For V959 Mon, the optical peak magnitude of 9.4 (unfiltered) was observed  $\sim 50$  days after the initial  $\gamma$ -ray detection, and we adopted an inferred peak of 5 magnitude (9).

Cyg (3). The  $\gamma$ -ray light curves of all four systems (Fig. 2) are similar, with 2- to 3-day-long peaks occurring 3 to 5 days after the initial LAT detections. The observed optical peak preceded the  $\gamma$ -ray peak by  $\sim 2$  days in V1324 Sco (13, 17) and  $\sim 6$  days in V339 Del (12, 13). Because the early optical light variations of the ejecta in novae are driven by line opacity changes in the ultraviolet during the expansion, the rise to peak optical brightness coincides with the maximum flux redistribution toward lower energies as the optically thick surface moves outward [see (18)]. The initial lack of detected  $\gamma$  rays could be because the ejecta are opaque and any  $>100$ -MeV emissions produced are absorbed by photon-atom interactions, with  $\gamma$  rays appearing only later when the density drops and the ejecta become transparent. The three novae were detected in  $\gamma$  rays during a time of high x-ray and ultraviolet/optical opacity. Coincidentally, the few days' delay of the  $\gamma$ -ray peak relative to the optical peak was also observed in V407 Cyg, but this may instead signal interactions with its red giant companion (below).

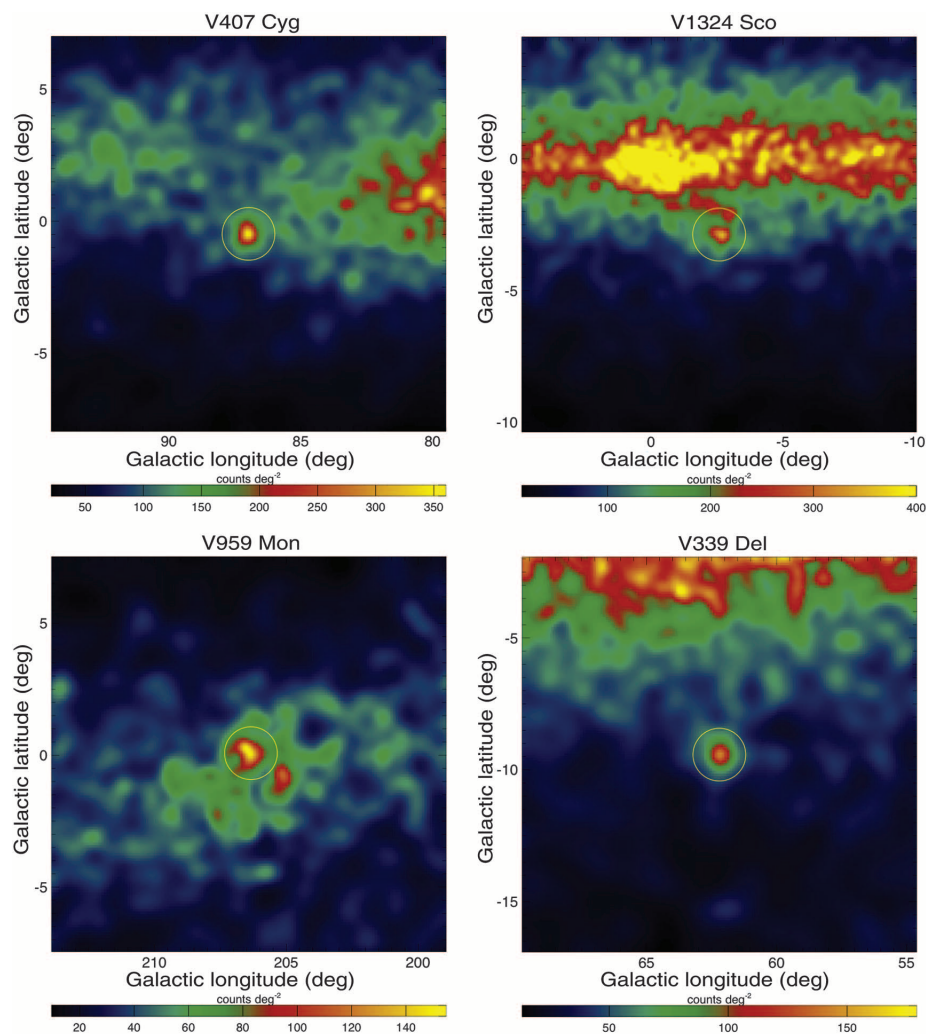
In compact classical nova binaries, typical companion separations are  $a \sim 10^{11}$  cm [ $\sim 100$  times

larger in symbiotic systems (19)], and expansion velocities  $v_{\text{ej}}$  at early times are many 100s to  $\geq 1000$  km  $s^{-1}$ . Thus, the ejecta reach the companion on a time scale of  $t = 1000 (a/10^{11} \text{ cm}) (v_{\text{ej}}/1000 \text{ km s}^{-1})^{-1}$  s (i.e., on the order of an hour or less). Modeling of the optical line profiles indicates that the spatial distribution of the ejected gas is bipolar rather than spherical in all cases, with greater extension perpendicular to the orbital plane in V959 Mon (9, 20, 21). Also, narrow absorption and emission line structures seen in optical and ultraviolet line profiles later in the expansion may be evidence of hydrodynamical instabilities and multiple ejections that may lead to the formation of strong turbulence and internal shocks within the ejecta after the ignition of the thermonuclear runaway (22). A clue to the physical process that causes the  $\gamma$ -ray emission mechanism may be the similarity of the high-energy spectral characteristics of V1324 Sco, V959 Mon, and V339 Del. Their  $>100$ -MeV spectra are all soft and can be fit with single power laws [the spectrum  $N(E) \propto E^{-\Gamma}$ , where  $N$  is the number of photons and  $E$  is energy] with photon indices  $\Gamma = 2.1$  to 2.3, or

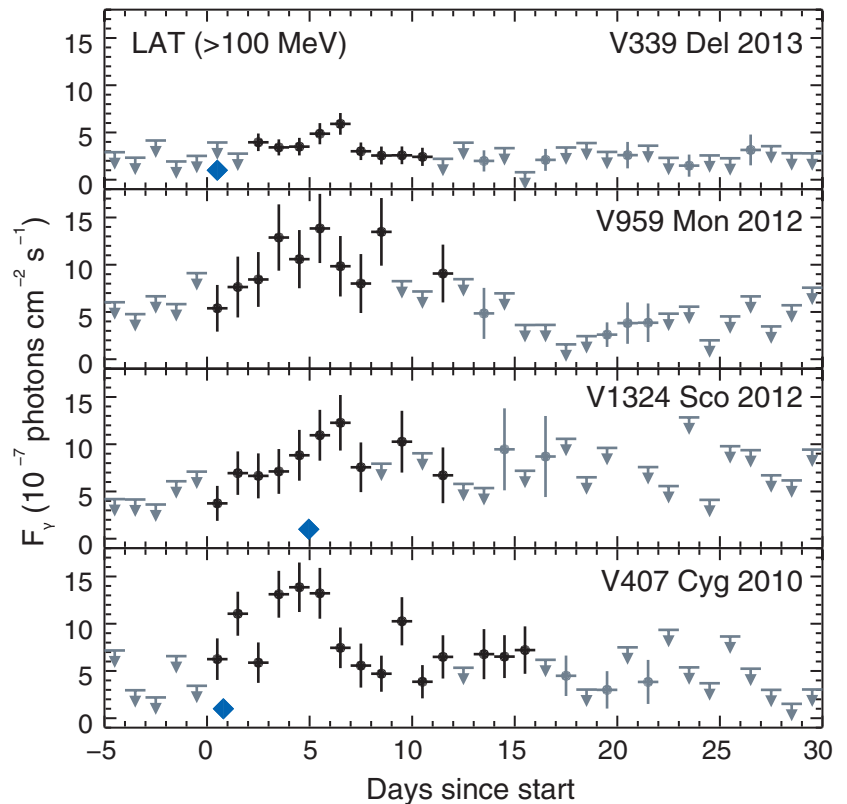
exponentially cutoff power laws [the spectrum  $N(E) \propto E^{-s} e^{-E/E_c}$ , where  $E_c$  is the cutoff energy] [see (13), table S1, and fig. S1]. The exponentially cutoff power-law fits to the LAT data were preferred over the power-law fits at the  $3.8 \sigma$  and  $3.4 \sigma$  level for V959 Mon and V339 Del, respectively, but provided an insignificant improvement ( $2.0 \sigma$ ) for V1324 Sco. Considering the uncertainties in the spectral fits, the three novae are similarly characterized by slopes  $s = 1.7$  to 1.8,  $E_c \sim 1$  to 4 GeV, and observed emission up to  $\sim 6$  to 10 GeV. The total durations of the observed  $\gamma$  rays were also similar, being detected for 17 to 27 days at  $>2 \sigma$  statistical significances in daily bins (Fig. 2 and Table 1). Because the LAT-observed properties are similar, it is likely that the  $\gamma$ -ray emission of these classical novae has a similar origin, involving interactions of the accelerated high-energy protons (hadronic scenario) or electrons (leptonic scenario) within the ejecta.

In the hadronic scenario, high-energy protons that interact with nuclei produce neutral pions ( $\pi^0$ ), which decay into two  $\gamma$  rays. For a representative hadronic model, we assume an exponentially cutoff power law distribution of

**Fig. 1. Fermi-LAT  $>100$ -MeV  $\gamma$ -ray counts maps of the four novae in Galactic coordinates centered on the optical positions over the full 17- to 27-day durations.** The maps used  $0.1^\circ$  by  $0.1^\circ$  pixels and were adaptively smoothed with a minimum number of 25 to 50 counts per kernel. Each nova (located at the centers of the yellow circles with  $1^\circ$  radius, which is the approximate LAT 95% containment at 1 GeV) is observed near the bright diffuse  $\gamma$ -ray emission in the Galactic plane, with V959 Mon in particular observed directly through the plane ( $0^\circ$  latitude).



**Fig. 2. Fermi-LAT 1-day binned light curves of the four  $\gamma$ -ray detected novae.** Vertical bars indicate  $1\sigma$  uncertainties for data points with  $>3\sigma$  (black) and 2 to  $3\sigma$  (gray) significances; otherwise,  $2\sigma$  upper limits are indicated with gray arrows. Start times  $t_s$  (from top to bottom panels) of 16 August 2013, 19 June 2012, 15 June 2012, and 10 March 2010 were defined as the day of the first  $\gamma$ -ray detection. In V339 Del, there was a  $2.4\sigma$  detection in 0.5-day binned data beginning 16.5 August (13), the epoch of the optical peak (blue diamond in each panel).



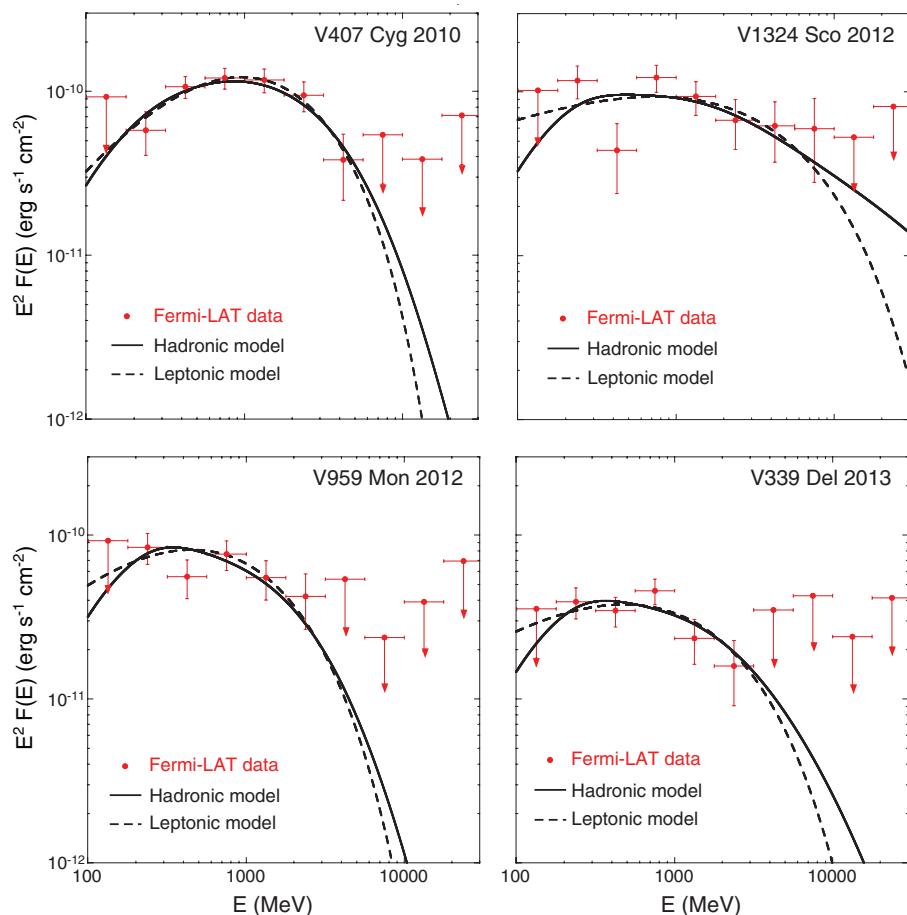
protons in the form  $N_p(p_p) = N_{p,0}(p_p c)^{-s_p} e^{-W_p/E_{cp}}$  (proton/GeV), where  $p_p$  and  $W_p$  are the momentum and the kinetic energy of protons, respectively;  $N_{p,0}$  the normalization;  $s_p$  the slope; and  $E_{cp}$  the cutoff energy. We fitted  $E_{cp}$  and  $s_p$  with the LAT spectra to obtain the best-fit  $\pi^0$  models (Fig. 3). The lower limits to the cutoff energies ( $\sim 3$  to  $30$  GeV) suggest proton acceleration up to near-TeV energies. The slopes of the best-fit models of the proton spectrum have large statistical uncertainties ( $\sim 0.8$ ) but, interestingly, are compatible with a value of 2 expected in the first-order Fermi acceleration process. To match the observed  $\gamma$ -ray variability time scale in such a process, a magnetic field  $B > 10^{-3}$  Gauss is required in a strong shock with  $v_{ej} = 2000$  km  $s^{-1}$  to accelerate particles to  $>1$  (10) GeV energies in  $\sim 0.2$  (2) days. Formally, the updated best-fit proton spectrum for the symbiotic nova V407 Cyg (3) is parameterized by  $s_p = 1.4^{+0.3}_{-0.4}$  GeV, but slopes of 2 to 2.2 are also viable at the 90% confidence level with  $E_{cp} = 10^{+1.0}_{-0.7}$  GeV (13) (fig. S3). Lower-confidence fits were also obtained for V959 Mon and V339 Del but, conversely, with smaller slopes and lower cutoff energies (13) (fig. S3). Assuming that the  $\gamma$ -ray flux is due to the interactions of high-energy protons with the nuclei in the ejecta, the best-fit parameters allow us to estimate the total energy in high-energy protons of  $\sim (3$  to  $17) \times 10^{42}$  ergs and to derive conversion efficiencies (i.e., the ratio of the total energy in high-energy protons to the kinetic energy of the ejecta) ranging from  $\sim 0.1$  to 3.7% for the classical novae and 6.6% for V407 Cyg.

In the leptonic case, accelerated electrons produce  $\gamma$  rays through a combination of inverse Compton scattering with low-energy photons and bremsstrahlung with atoms in the vicinity of the nova. For a leptonic model, we adopted a similar functional form for the distribution of the kinetic energy of high-energy electrons ( $W_e$ ) in the form  $N_e(W_e) = N_{e,0} W_e^{-s_e} e^{-W_e/E_{ce}}$  (electron/GeV) and fitted the normalization  $N_{e,0}$ , slope  $s_e$ , and cutoff energy  $E_{ce}$  to the LAT data for each nova (Fig. 3). The  $\gamma$ -ray luminosity of the calculated bremsstrahlung emission is  $<20\%$  of the total  $\gamma$ -ray luminosity for all the novae (13). The best-fit parameters of the high-energy electron spectra for the three classical novae are similar within their confidence regions (13), with  $E_{ce}$  constrained to lie between 2 and 30 GeV and poorly constrained slopes. These models are statistically indistinguishable from the  $\pi^0$  model. As in the hadronic model, the spectral parameters of the classical novae differ from those for V407 Cyg (mainly due to the lowest-energy  $\sim 200$ - to  $300$ -MeV bin detected in its LAT spectrum), where the best-fit slope is negative (i.e., a positive index of the power law) and  $E_{ce} = 1.78 \pm 0.05$  GeV. The best-fit parameters for the leptonic scenario, where high-energy electrons interact primarily with the photons emitted by the nova photosphere (23), lead to total energies of  $\sim (6$  to  $13) \times 10^{41}$  ergs in high-energy electrons and conversion efficiencies of  $\sim 0.1$  to 0.3% for the classical novae and 0.6% for the symbiotic system.

Detection of classical novae in  $\gamma$  rays was deemed unlikely in the past (3). The only nova

previously detected in  $\gamma$  rays, the aforementioned V407 Cyg, was a rare symbiotic and likely recurrent [only 10 recurrent novae are known, of which 4 are symbiotic types (24)]. In the symbiotic novae, conditions are conducive for high-energy particle acceleration as the portion of the ejecta moving into the wind in the direction of the dense medium provided by the red giant companion decelerates within a few days. The  $\gamma$  rays peak early, when the efficiency for hadron and lepton acceleration is presumably favorable, with the red giant wind playing a key role in the  $\gamma$ -ray production (2, 23). In contrast, the main-sequence star companions in the classical novae do not provide similarly dense target material; hence, it is likely that other dissipative processes are involved in particle acceleration and generation of the observed  $\gamma$  rays.

Because the  $\gamma$ -ray properties of the novae detected so far by the Fermi-LAT appear similar to one another, and their underlying properties are unremarkable, it appears that all novae can be considered to be candidate  $\gamma$ -ray emitters. Their detection by the LAT may imply close proximity and that other optical novae not yet detected with the LAT [e.g., (25)] are more distant and have fainter optical peaks [without considering extinction uncertainties (26)]. Indeed, all the LAT-detected novae have estimated distances of  $\lesssim 4$  to 5 kpc (Table 1). Despite systematic uncertainties in the adopted distances, it is interesting that the inferred mean  $\gamma$ -ray luminosities and total emitted energies of the novae span a small range  $\sim (3$  to  $4) \times 10^{35}$  ergs  $s^{-1}$  and  $\sim (6$  to  $7) \times 10^{41}$  ergs, respectively, except for the  $\sim 2$  times



**Fig. 3.** Fermi-LAT >100-MeV average  $\gamma$ -ray spectra of the four novae over the full 17- to 27-day durations. Vertical bars indicate  $1\sigma$  uncertainties for data points with significances  $>2\sigma$  otherwise, arrows indicate  $2\sigma$  limits. The best-fit hadronic and leptonic model curves are overlaid.

greater values for V1324 Sco, whose distance is highly uncertain.

The rate of novae in the Milky Way is highly uncertain, but considering a plausible range of  $\sim 20$  to 50 per year (27) and reasonable spatial distributions in the Galactic bulge and disk (28), our estimate is 1 to 4 per year at  $\leq 4$ - to 5-kpc distances. The  $\gamma$ -ray detection rate of novae averages roughly one per year over the time span of these observations ( $\sim 5$  years), consistent with the lower end of this extrapolation.

Although the  $\gamma$ -ray properties of the LAT-detected novae are similar, we emphasize the small and subtle differences that imply different emission mechanisms—e.g., the spectral shape of V407 Cyg compared with the three classical novae as well as the apparent higher energy extension of the V1324 Sco spectrum. Among the classical novae detected so far, they also appear different optically. The  $\gamma$ -ray emission mechanism and high-energy particle acceleration processes associated with the novae could depend on the particular system properties that remain to be investigated, such as the white dwarf mass, which determines the explosion energetics (ejected mass and expansion velocity), and the mass transfer dictated by the companion mass and separation.

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#### SUPPLEMENTARY MATERIALS

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## QUANTITATIVE SOCIAL SCIENCE

# A network framework of cultural history

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The emergent processes driving cultural history are a product of complex interactions among large numbers of individuals, determined by difficult-to-quantify historical conditions. To characterize these processes, we have reconstructed aggregate intellectual mobility over two millennia through the birth and death locations of more than 150,000 notable individuals. The tools of network and complexity theory were then used to identify characteristic statistical patterns and determine the cultural and historical relevance of deviations. The resulting network of locations provides a macroscopic perspective of cultural history, which helps us to retrace cultural narratives of Europe and North America using large-scale visualization and quantitative dynamical tools and to derive historical trends of cultural centers beyond the scope of specific events or narrow time intervals.

Quantifying historical developments is crucial to understanding a large variety of complex processes from population dynamics to disease spreading, conflicts, and urban evolution. However, in historical research there is an inherent tension (*1, 2*) between qualitative analyses of individual historical accounts and quantitative approaches aiming to measure and model more general patterns. We believe that these approaches are complementary: We need quantitative methods to identify statistical regularities, as well as qualitative approaches to

explain the impact of local deviations from the uncovered general patterns. We have therefore developed a data-driven macroscopic perspective that offers a combination of both approaches.

We collected data from Freebase.com (*3*), the General Artist Lexicon (AKL) (*4–6*), and the Getty Union List of Artist Names (ULAN) (*7*), representing spatiotemporal birth and death information of notable individuals, spanning a time period of more than two millennia. The data sets are included in the supplementary materials (SM), accompanied by an explanation of their nature and data preparation (*8*) (tables S1 and S2). Potential sources of bias are addressed in the SM, including biographical, temporal, and spatial coverage; curated versus crowd-sourced data; increasing numbers of individuals who are still alive; place aggregation; location name changes and spelling variants; and effects of data set language. Most important, compared with contemporary worldwide migration flux (*9*), our data sets focus on birth-to-death migration within and out of Europe and North America (see fig. S1). Notability of individuals, simply defined as the curatorial decision of inclusion in the respective data set, differs slightly between the more

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